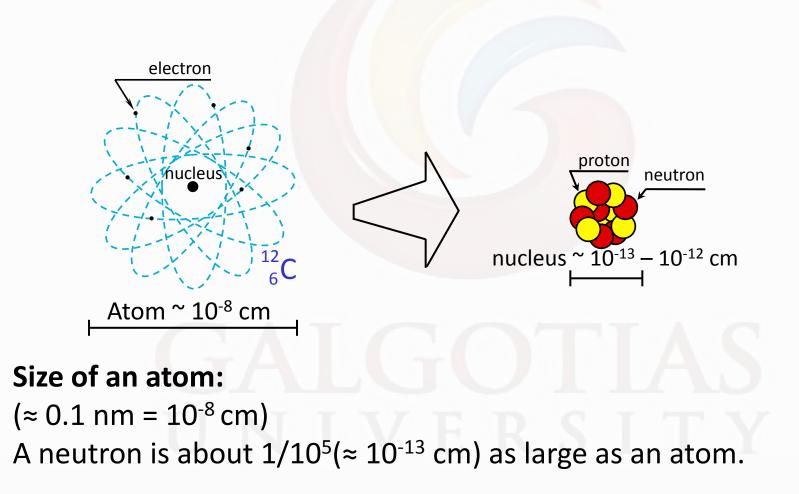


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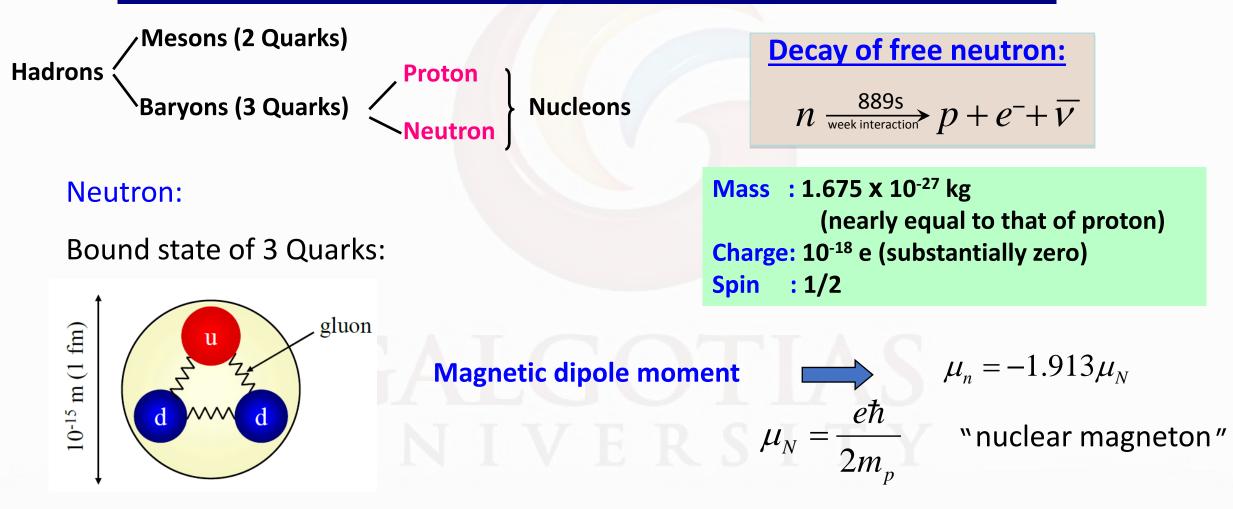
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Particle Nature of Neutron



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Course Code : MSCP6015 Course Name: Nuclear and Particle Physics 1 Particle as a Wave $E = k_{\rm B}T$ (From classical physics) To calculate the interference $E = \hbar \omega$ (Plank's hypothesis) effect during the scattering $p = \frac{E}{c} = \frac{\hbar\omega}{c}$ (Plank-Einstein relation) $\vec{p} = h\vec{k}$ $p = \frac{h}{\lambda}$ (de Broglie hypothesis) process, a particle has to be described as a matter wave. (here, ω =ck and the fact that the momentum and wave vector point in the same direction) Kinetic energy of neutron: $E_n = \frac{p^2}{2m_n} = \frac{h^2}{2m_n\lambda^2} \equiv k_B T_{eq}$ Wave length of neutron: $\lambda = \frac{h}{\sqrt{2m_nk_BT}}$

Name of the Faculty: Dr. Shyamal Kumar Kundu

Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

Scattering

Internal structure:

what are the relevant building blocks (atoms, molecules, colloidal particles, ...) and how are they arranged?

Microscopic dynamics:

How do these building blocks move (atomic movements) and what are their internal degrees of freedom?

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Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

Scattering

For magnetic systems, we need to know the arrangement of the microscopic magnetic moments due to spin and orbital angular momentum and their excitation spectra.

The macroscopic response and transport properties, such as specific heat, thermal conductivity, elasticity, viscosity, susceptibility, magnetization etc., which are the quantities of interest for applications, result from the microscopic structure and dynamics.

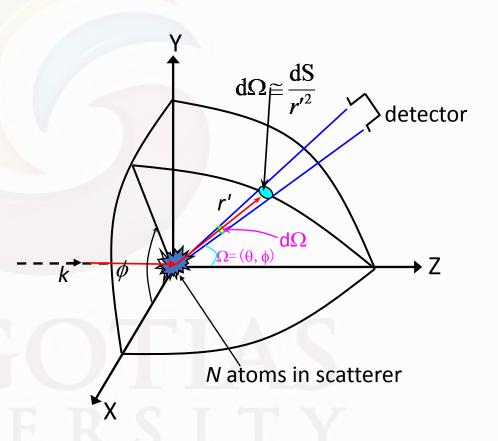
Course Code : MSCP6015

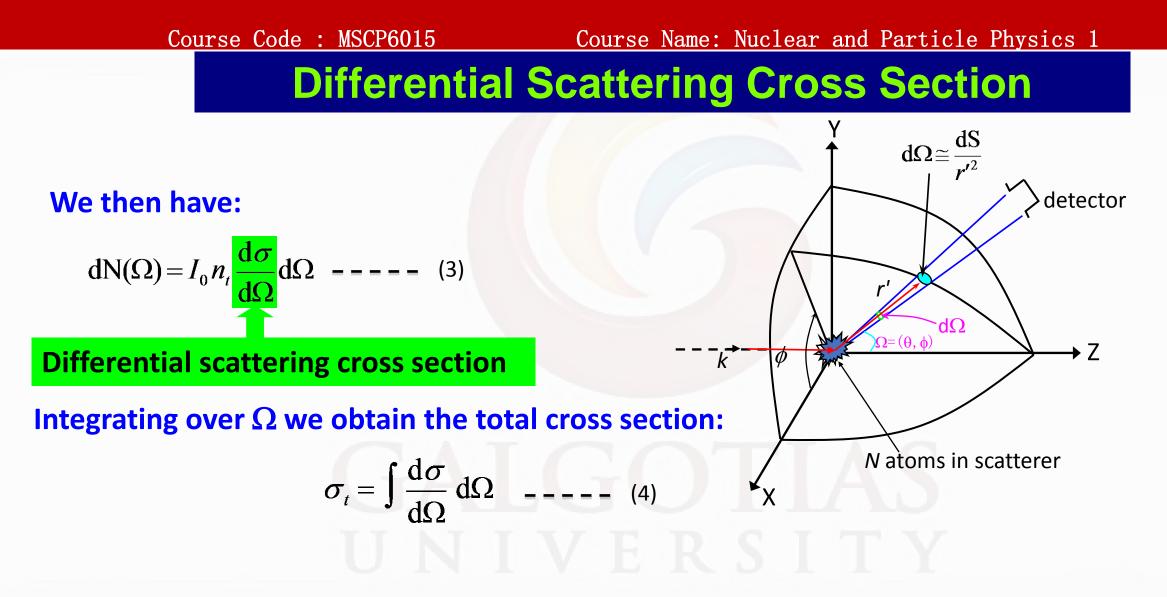
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Differential Scattering Cross Section

The number of particles $\Delta N(\Omega)$ per unit time per unit target volume that have undergone a collision and are recorded by the detector is proportional to i) The flux I_0 of incident particles

- ii) The density n_{t_i} of the target particle
- iii) The solid angle $d\Omega$ the detector subtends as seen from the target





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Course Code : MSCP6015

→Z

 $d\Omega \approx \frac{dS}{r^2}$

 \overline{k}

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Scattering From a Static System

A quantum mechanical wave function can be A describe by:

 $\psi(\mathbf{r},t) = A \exp(i(\varpi t - \mathbf{k} \cdot \mathbf{r})) \qquad \text{----} \qquad (5)$

where, $\boldsymbol{\varpi} = 2\pi f$, $k = 2\pi / \lambda$

N atoms in scatterer The probability of the observation of a particle can be described by:

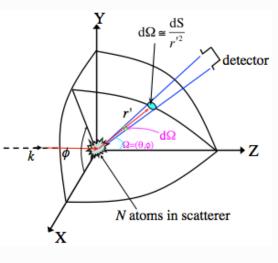
$$I = |\psi(\mathbf{r}, t)|^2 \qquad \dots \qquad (6)$$
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Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

Scattering From a Static System



The wave scattered by a single scattering center is given by:

$$\psi(\mathbf{r}',t) = \frac{\exp(-i(k|\mathbf{r}'-\mathbf{r}_{j}|))}{|\mathbf{r}'-\mathbf{r}_{j}|} b_{j}\psi(\mathbf{r},t)$$
 ---- (7)

Here,

 $r' \Rightarrow$ point of observation of the scattered wave

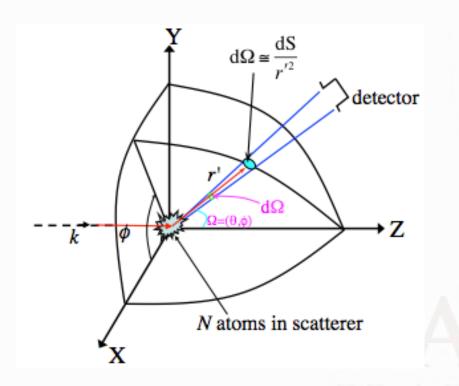
 $r_i \Rightarrow$ position of the scattering center

 $b_i \Rightarrow$ property of this scattering center, the scattering length

 $exp(-i(k|r'-r_j|)) \Rightarrow$ additional oscillations between the scattering process and detection

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Course Code : MSCP6015 Course Name: Nuclear and Particle Physics 1 Scattering From a Static System

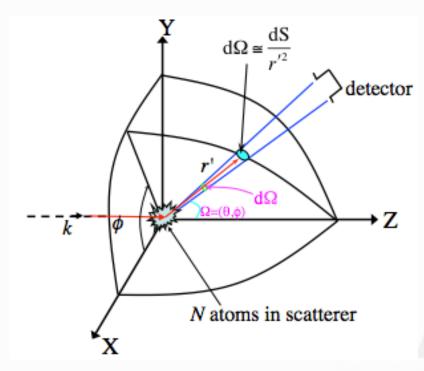


Now we assumed that the scattering centres are fixed points in space described by coordinates $r_j(t)$ =const. = r_j . In this case the scattering is elastic so that before and after the scattering process energy will not change.

Name of the Faculty: Dr. Shyamal Kumar Kundu

 Course Code : MSCP6015
 Course Name: Nuclear and Particle Physics 1

 Scattering From a Static System

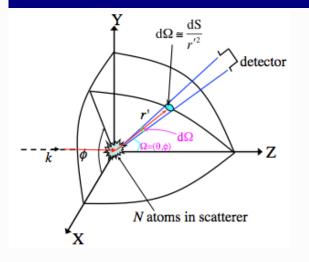


Without loss of generality one can define the direction of the incident beam as the z direction: k=(0,0,k). By using eq. 5 and eq. 7, the wave scattered by particle j into the detector can be expressed as

$$\psi(\mathbf{r}',t) = \frac{\exp(-i(k|\mathbf{r}'-\mathbf{r}_j|))}{|\mathbf{r}'-\mathbf{r}_j|} Ab_j \exp(i(\varpi t - \mathbf{k} \cdot \mathbf{r})) \quad ---- \quad (8)$$

Name of the Faculty: Dr. Shyamal Kumar Kundu

Course Code : MSCP6015 Course Name: Nuclear and Particle Physics 1 Scattering From a Static System



Now, we assumed that SDD, R, is large compared to the size of the sample (Fraunhofer diffraction). Then an expression to 1st order in the scatterer ' s position yields:

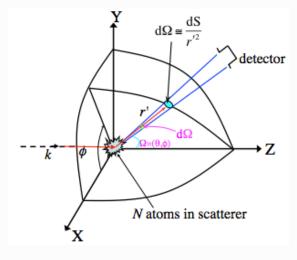
$$\psi(\mathbf{r}',t) = \frac{A}{R} e^{i \, \omega t} e^{i k (R+L)} b_j \exp(i\mathbf{Q} \cdot \mathbf{r}_j) \quad ---- \quad (9)$$

Then the intensity:
$$I = \psi \cdot \psi^* = \frac{A^2}{R^2} \left| \sum_{j=1}^N b_j \exp(iQ \cdot r_j) \right|^2$$
 ---- (10)

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 Course Code : MSCP6015
 Course Name: Nuclear and Particle Physics 1

Scattering From a Static System



The differential scattering cross section can be written as:

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \left\langle \left| \sum_{j=1}^{N} b_j \exp(\mathrm{i}\mathbf{Q} \cdot \mathbf{r}_j) \right|^2 \right\rangle = \left\langle \left| A(Q) \right|^2 \right\rangle \quad ---- \quad (11)$$

Where, $A(Q) = \sum_{i} b_{j} \exp(iQ \cdot r_{j})$ ---- (11a)

is the superposition of the individual scattering amplitudes.

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Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

Coherent and Incoherent Scattering

Since, the individual scattering lengths are independent, then, we can write:

$$\langle b_i b_j \rangle = \langle b \rangle^2 + (\langle b^2 \rangle - \langle b \rangle^2) \delta_{ij}$$
 ----- (12)

Where, δ_{ij} is the delta function

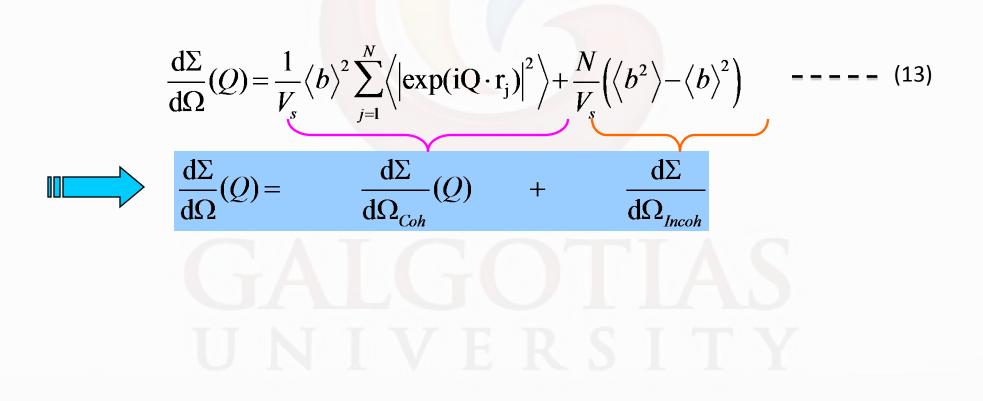
The macroscopic differential scattering cross section can be written as: $\frac{d\Sigma}{d\Omega} = \frac{1}{V_s} \frac{d\sigma}{d\Omega}$

Name of the Faculty: Dr. Shyamal Kumar Kundu

Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

Coherent and Incoherent Scattering

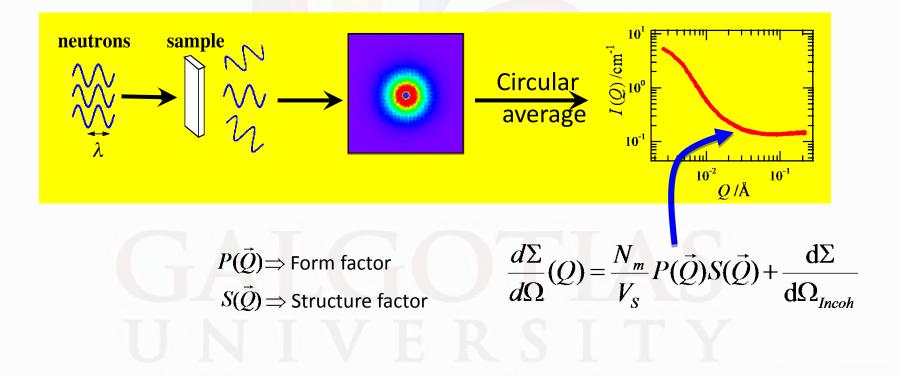


Name of the Faculty: Dr. Shyamal Kumar Kundu

Course Code : MSCP6015

Course Name: Nuclear and Particle Physics 1

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